Intercalation of Trioxatriangulenium Ion in DNA: Binding, Electron Transfer, X-ray Crystallography, and Electronic Structure

Johannes Reynisson†,‡ Gary B. Schuster,*,† Sheldon B. Howerton,‡ Loren Dean Williams,*,† Robert N. Barnett,§ Charles L. Cleveland,§ Uzi Landman,*,§ Niels Harrit,‡ and Jonathan B. Chaires†

Contribution from the Schools of Chemistry and Biochemistry and of Physics, Georgia Institute of Technology, Atlanta, Georgia 30332, Nano-science Center, University of Copenhagen, H.C., Ørsted Institute DK-2100, Copenhagen, Denmark, and Department of Biochemistry, University of Mississippi Medical Center, Jackson, Mississippi 39216

Received August 23, 2002; E-mail: gary.schuster@cos.gatech.edu

Abstract: Trioxatriangulenium ion (TOTA+) is a flat, somewhat hydrophobic compound that has a low-energy unoccupied molecular orbital. It binds to duplex DNA by intercalation with a preference for G-C base pairs. Irradiation of intercalated TOTA+ causes charge (radical cation) injection that results in strand cleavage (after piperidine treatment) primarily at GG steps. The X-ray crystal structure of TOTA+ intercalated in the hexameric duplex d[GATCG]6 described here reveals that intercalation of TOTA+ results in an unusually large extension of the helical rise of the DNA and that the orientation of TOTA+ is sensitive to hydrogen-bonding interactions with backbone atoms of the DNA. Electronic structure calculations reveal no meaningful charge transfer from DNA to TOTA+ because the lowest unoccupied molecular orbital of TOTA+, (LUMO), falls in the gap between the highest occupied molecular orbital, (HOMO)D, and the (LUMO)B of the DNA bases. These calculations reveal the importance of backbone, water, and counterion interactions, which shift the energy levels of the bases and the intercalated TOTA+ orbitals significantly. The calculations also show that the inserted TOTA+ strongly polarizes the intercalation cavity where a sheet of excess electron density surrounds the TOTA+.

Introduction

Interest in understanding the association of small molecules with duplex DNA is driven by recognition that the precise function of numerous synthetic and natural products that affect DNA chemically or structurally is determined by the mode of binding.1,2 Intercalation and groove binding are the two most common forms that this association takes. Groove binding is typically observed for arc-shaped, nonplanar compounds that have functional groups able to associate with the hydrogen-bond donors and acceptors found in the major and minor grooves. Binding by intercalation is characteristic of planar, hydrophobic compounds that effectively fill the space produced between base pairs formed when the helix is locally elongated and partially unwound. Intercalated compounds typically interact with adjacent base pairs primarily through van der Waal forces and electrostatic stabilization. Consequently, it appears that polarizable compounds intercalate preferentially at sites containing G-C base pairs since their dipole moment is greater than that of their A-T counterparts.3

Extensive investigation of the electronic interaction of compounds intercalated in DNA follows from the important biological and medical roles played by intercalators. Additional interest in investigation of intercalation and stacking was stimulated by the report of accelerated electron transfer between metal complexes mediated by coupling of donor and acceptor states of these compounds with electronic states of the DNA bases.4 The decisive role of electronic coupling is highlighted by the claimed contrasting behavior of (i) metallointercalators, which elicit the “wirelike” behavior of DNA,5 (ii) intercalated, tethered ethidium, which must reorient to enable wirelike behavior,6 and (iii) electron transfer between modified nucleic acid bases, which reveals the invalidity of wirelike behavior.7

It seemed that TOTA+ is an ideal compound for examination of the structural, electronic, and chemical effects of an intercalator having a relatively low-energy unfilled orbital. Trioxatriangulenium carbocation (TOTA+, Figure 1)8 is a stable, planar9 compound with a relatively low-energy LUMO.10,11


1 School of Chemistry and Biochemistry, Georgia Institute of Technology.  
2 University of Copenhagen.  
3 School of Physics, Georgia Institute of Technology.  
4 School of Mississippi Medical Center.

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Optical excitation of TOTA\(^+\) with visible light in aqueous solution forms excited states capable of the one-electron oxidation of guanosine or adenosine.\(^9\)

We report here an investigation of the binding of TOTA\(^+\) to DNA oligomers in solution and the photophysical and photochemical reactions of bound TOTA\(^+\). Irradiation of bound TOTA\(^+\) results in damage at GG steps that is revealed as strand cleavage following treatment with piperidine. This cleavage indicates that light-induced electron transfer from the DNA to the TOTA\(^+\) has occurred. We obtained a high-resolution X-ray diffraction of TOTA\(^+\) intercalated in the self-complementary hexameric DNA duplex (d(CGATCG))\(_2\). The structure of the DNA:TOTA\(^+\) complex allows us to (i) confirm that TOTA\(^+\) is an intercalator, (ii) assess the interactions that govern binding orientation and sequence selectivity, (iii) determine the structural consequences to DNA of intercalation of a carbocationic compound that localizes charge near the center of the base pair stack, and (iv) test and extend previous hypotheses on relationships between intercalator structure and DNA distortion.\(^10\)

We also report the results of first-principles quantum mechanical calculations on TOTA\(^+\) intercalated in DNA that assess the nature of electronic interactions with adjacent groups of the DNA. To explore the electronic structure and binding characteristics of TOTA\(^+\), we performed a series of calculations with the crystallographically determined atomic positions as part of the input; and (ii) we employed quantum mechanical calculations for evaluation of the electronic energy level spectrum of the intercalated DNA molecule, using configurations determined by the MD simulations, for investigation of the bonding mechanism of the intercalated TOTA\(^+\) cation to the DNA. These calculations reveal that, despite its low-lying unoccupied orbital, the interactions of TOTA\(^+\) with DNA are controlled primarily by hydrophobic, hydrogen-bonding, and electrostatic interactions. There is no significant charge transfer from the DNA to TOTA\(^+\) in the ground state.

Materials and Methods

4,8,12-Trioxa-4,8,12,12c-tetrahydrodibenzo[cd,mn]pyrenylum carboxation (TOTA\(^+\)) tetrafluoroborate was synthesized according to the procedure published by Martin and Smith.\(^8\)

**Determination of the Binding Constant.** A sodium phosphate buffer (10 mM) solution of TOTA\(^+\) (pH = 7.0, 25 \(\mu\)M in a 1.0 cm cuvette) was titrated with a concentrated calf thymus DNA (CT-DNA) solution. The absorbance of the solution at 330 nm was measured initially and after each addition of CT-DNA. The change in absorbance and the extinction coefficients for TOTA\(^+\) in buffer and CT-DNA were used to calculate \(c\), the fraction of DNA-bound TOTA\(^+\).\(^14\)

**Fluorescence Quenching.** Solutions of 25 \(\mu\)M TOTA\(^+\) and CT-DNA were prepared in phosphate buffer solution (pH = 7.0) and transferred to a standard 1 cm optical quartz cell. The emission spectrum from 475 to 650 nm was obtained at room temperature on a Spex 1681 Fluorolog emission spectrometer in the normal 90° configuration with excitation at 450 nm.

**Solution Viscosity Measurements.** Viscosity experiments used a Cannon–Ubbelohde semimicro dilution viscometer immersed in a thermostated water bath maintained at 33.3 ± 0.2°C. Small volumes of a concentrated TOTA\(^+\) stock solution (5 \(\times\) 10\(^{-4}\) M) were added to sonicated CT-DNA samples (average 800 base pairs long, 2.4 \(\times\) 10\(^{-4}\) M in base pairs) in the viscometer. These solutions were mixed by bubbling them with air. The TOTA\(^+\) and DNA solutions were buffered with piperazine-N,N′-bis(2-ethanesulfonic acid) at pH 7.0. Relative viscosities for DNA in either the presence or absence of TOTA\(^+\) were calculated from the relation \(η = (t - t_0)/t_0\), where \(t\) is the observed flow time of the DNA-containing solution and \(t_0\) is the flow time of the buffer alone. Viscosity data were plotted as \((η/η_0)\) versus the binding ratio \(r\), according to the theory of Cohen and Eisenberg.\(^15\)

**Induced Circular Dichroic (CD) Spectroscopy.** CD spectra of TOTA\(^+\) (30 \(\mu\)M in 10 mM sodium phosphate, pH 7) and sonicated CT-DNA (40 \(\mu\)M, base pairs) were recorded on a Jasco J-270 spectrometer in a 10 cm path length cylindrical quartz cell from 200 to 550 nm.

**Melting Temperatures (\(T_m\)).** Melting curves were obtained for TOTA\(^+\)-containing phosphate buffer solutions of two complimentary 12-mer oligonucleotides d(5′-AAA GGT ACG GGG-3′) and d(5′-CGC GTT ACC TTT-3′), 5 \(\mu\)M each in strand, by monitoring the absorbance at 260 nm on Cary 1E UV-visible spectrophotometer as the temperature was ramped from 20 to 90°C at the rate of 0.5°C/min. Melting curves obtained from heating and cooling ramps are essentially identical.

**Competition Dialysis.** Multiple disposable dialysis units (Spectrum Laboratories, Inc.), each containing 0.5 mL of a 75 \(\mu\)M solution of the DNA sample listed in Table 1, were placed in a beaker containing 200 mL of a 1 \(\mu\)M TOTA\(^+\) solution. The system was allowed to equilibrate for 24 h with continuous stirring. After equilibrium was reached, the absorbance of TOTA\(^+\) bound to the DNA was measured by extraction into an SDS solution. Since all of the DNA samples are in equilibrium with the same TOTA\(^+\) concentration, the amount of bound ligand is directly proportional to the association constant for binding of TOTA\(^+\).\(^16\)

**Cleavage Analysis.** The 12-mer duplex d(5′-AAA GGT ACG GGG-3′) was radiolabeled at the 5′-terminus by use of (γ-\(^32\)P)ATP and bacterial T4 polynucleotide kinase\(^17\) and purified by 20% PAGE. Samples for irradiation were prepared by hybridizing a mixture of cold and radiolabeled oligonucleotide (5 \(\mu\)M) with 5 \(\mu\)M the complementary strand and 10 \(\mu\)M TOTA\(^+\) in a phosphate buffer solution. The samples were irradiated in microcentrifuge tubes in a Rayonet photoreactor containing eight 350 nm lamps. After irradiation, the samples were precipitated with ethanol, dried, and then treated with piperidine (1 M in 0.1 M NaOH) for 3 h. The products were analyzed by polyacrylamide gel electrophoresis.

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Table 1. Nucleic Acid Conformations and Samples Used in Competition Dialysis Experiments

<table>
<thead>
<tr>
<th>conformation</th>
<th>DNA oligonucleotide</th>
<th>λ² (nm)</th>
<th>ε (M⁻¹ cm⁻¹)</th>
<th>monomeric unit</th>
</tr>
</thead>
<tbody>
<tr>
<td>single-strand purine</td>
<td>poly(dA)</td>
<td>257</td>
<td>8600</td>
<td>nucleotide</td>
</tr>
<tr>
<td>single-strand pyrimidine</td>
<td>poly(dT)</td>
<td>264</td>
<td>8520</td>
<td>nucleotide</td>
</tr>
<tr>
<td>duplex DNA</td>
<td>C. perfringens (31% GC)</td>
<td>260</td>
<td>12460</td>
<td>base pair</td>
</tr>
<tr>
<td></td>
<td>calf thymus (42% GC)</td>
<td>260</td>
<td>12824</td>
<td>base pair</td>
</tr>
<tr>
<td></td>
<td>M. lysodeikticus (72% GC)</td>
<td>260</td>
<td>13846</td>
<td>base pair</td>
</tr>
<tr>
<td></td>
<td>poly(dA)poly(dT)</td>
<td>260</td>
<td>12000</td>
<td>base pair</td>
</tr>
<tr>
<td></td>
<td>[poly(dA)⋅poly(dT)]</td>
<td>262</td>
<td>13200</td>
<td>base pair</td>
</tr>
<tr>
<td></td>
<td>[poly(dG⋅dC)]</td>
<td>254</td>
<td>16600</td>
<td>base pair</td>
</tr>
<tr>
<td>DNA-RNA hybrid</td>
<td>poly(dA)poly(dT)</td>
<td>260</td>
<td>12460</td>
<td>base pair</td>
</tr>
<tr>
<td></td>
<td>poly(dA)poly(dT)</td>
<td>260</td>
<td>14280</td>
<td>base pair</td>
</tr>
<tr>
<td></td>
<td>Z DNA⁺ brominated [poly(dG⋅dC)]</td>
<td>254</td>
<td>16060</td>
<td>base pair</td>
</tr>
<tr>
<td></td>
<td>triplex DNA</td>
<td>260</td>
<td>17200</td>
<td>base triplet</td>
</tr>
<tr>
<td></td>
<td>tetraplex DNA 1</td>
<td>260</td>
<td>39267</td>
<td>base tetrad</td>
</tr>
</tbody>
</table>

* λ is the wavelength of maximum absorbance. * ε is the molar extinction coefficient at λ, expressed in terms of the monomeric unit. * Prepared according to ref 60.

MD Simulations. The classical MD simulations employed the Amber 6 program suite with the AMBER/96 force field (25) supplemented by partial charges determined by us for TOTA⁺; see the Supporting Information. In these simulations, the computational cell contained the [d(CGATCG)]₂ duplex (including the sugar–phosphate backbone), with TOTA⁺ intercalated between the two CG steps (i.e., [dT/TOTA⁺/GATC/TOTA⁺/G]). The atoms of the DNA bases, the backbone, and the TOTA⁺ were kept (except for the hydrogen atoms) in their crystallographically determined positions. In addition, the computational cell contained eight Na⁺ ions and 1610 water molecules, which, together with the DNA hydrogen atoms, were treated dynamically; in constant pressure MD simulations of the system at atmospheric pressure and a temperature of 300 K (STP conditions), the average dimensions of the Cartesian calculational cell were 36.4 Å × 35.7 Å × 41.5 Å, with the larger dimension in the direction of the DNA helix axis. In the MD simulations the equations of motion were integrated by use of the Verlet algorithm with a time step of 1 fs.

Quantum Calculations. First-principles quantum mechanical calculations were performed by the Born–Oppenheimer (BO) local spin indicated that the sequence of the search model was incorrect. The sequence of the model was corrected and the two TOTA⁺ molecules were manually rotated to improve the fit into electron density. The simulated-anneling routine in CNS was performed to help remove bias from the manipulated model. A magnesium ion was added to the model, where it was clearly indicated in the Fourier electron density by its octahedral geometry and characteristic water–magnesium distances. Water molecules were added iteratively to peaks of corresponding sum and difference density followed by refinement and phase recalculation. After the model was completed, the R-free was dropped and all data were used to refine the final model. Annealed omit maps were calculated from a model excluding the two TOTA⁺ molecules. All atoms of both DNA strands and both TOTA⁺ molecules fit cleanly into the 2Fo − Fc and Fc − Fo omit electron density. The geometry of the refined complexes was characterized with the program Curves. (24)

References:

density (LSD) molecular dynamics (MD) method [BO-LSD-MD],\textsuperscript{27} with gradient corrections (generalized gradient approximation, GGA) for the exchange-correlation functional.\textsuperscript{28} The Kohn–Sham equations were solved in conjunction with norm-conserving nonlocal pseudopotentials for the valence electrons\textsuperscript{29} and a plane-wave basis (i.e., no atom-centered basis functions were used) with a high kinetic energy cutoff of 845 eV; a calculation for the three-base-pair duplex, an intercalated TOTA\textsuperscript{+}, 64 H\textsubscript{2}O molecules, and three Na atoms involves 414 ions and 1248 valence electrons. The algorithm for solving the density-functional Kohn–Sham equations uses a Fermi distribution function for the electrons, which is a very effective way of dealing with degenerate or near-degenerate energy levels.\textsuperscript{27} The Fermi temperature used is rather low, i.e., 0.01 eV/$k_B$, which ensures that the Fermi function is operative only on the nearly degenerate levels at the top of the energy level spectrum and not anywhere else (where the spectral gaps are larger). The BO-LSD-MD method is particularly suitable for calculations of charged systems or systems where higher multipole moments may develop, since no periodic replication of the ions is imposed; that is, no supercells are employed.\textsuperscript{27} Earlier simulation of hydrated DNA sequences also used the BO-LSD-MD method.\textsuperscript{30}

Results and Discussion

(I) Binding and Light-Induced Reactions of TOTA\textsuperscript{+} with DNA Oligomers in Solution: (I) Association of TOTA\textsuperscript{+} with DNA. The association of a small molecule with DNA often causes measurable changes to its electronic absorption spectrum that can indicate the mode and strength of binding.\textsuperscript{31} We examined the absorption and fluorescence spectra of TOTA\textsuperscript{+} in the presence of CT-DNA. The three lowest energy absorption bands for TOTA\textsuperscript{+} in aqueous buffer solution exhibit a pronounced hypochromic effect as well as a bathochromic shift (540–830 cm$^{-1}$) when CT-DNA is added to the mixture; see Figure 2. The bathochromic effect is most pronounced for the lower energy transition (\(\lambda_{max} = 480 \text{ nm}\)), and the hypochromicity is greatest for the higher energy transition (\(\lambda_{max} = 330 \text{ nm}\)).

The equilibrium binding constant for TOTA\textsuperscript{+} to DNA in 10 mM sodium phosphate buffer solution at pH 7 was estimated by monitoring the change of its absorption spectrum upon incremental addition of known amounts of CT-DNA. The absorbance of the band with \(\lambda_{max} = 330 \text{ nm}\) decreased and shifted to lower energy (red) after each addition. Scatchard analysis\textsuperscript{32} of these data yields a linear fit, consistent with a single binding mode of TOTA\textsuperscript{+} to DNA. The estimated binding constant is \((4.1 \pm 0.3) \times 10^4 \text{ M}^{-1}\), and the extinction coefficient of the TOTA\textsuperscript{+}-DNA complex is estimated to be 12 200 M$^{-1}$ cm$^{-1}$ at 336 nm. It is difficult to measure the extinction coefficients for the two other transitions of TOTA\textsuperscript{+}, which are in the visible range, due to their overlap.

The circular dichroic (CD) spectrum of TOTA\textsuperscript{+} bound to DNA provided additional information about the nature of this complex. Binding of an achiral molecule within a chiral environment, such as DNA, can lead to induced optical activity of the bound species.\textsuperscript{33} This effect has been observed for a variety of DNA intercalators.\textsuperscript{34} Since free TOTA\textsuperscript{+} is achiral, it has no CD spectrum; however, it exhibits induced CD bands upon binding to CT-DNA as is shown in Figure 3. These bands are relatively weak, which suggests that binding of TOTA\textsuperscript{+} occurs by intercalation.\textsuperscript{35,36}

The electronic interaction of TOTA\textsuperscript{+} bound to DNA can be assessed by fluorescence spectroscopy. The fluorescence of TOTA\textsuperscript{+} is quenched by the four individual DNA nucleosides. The Stern–Volmer plots are linear and reveal apparent quenching constants \((k_q)\) of \(1.5 \times 10^{10} \text{ M}^{-1} \text{ s}^{-1}\) for 2’-deoxyguanosine and 2’-deoxyadenosine, which is consistent with values reported previously.\textsuperscript{12} The fluorescence quenching of TOTA\textsuperscript{+} by the pyrimidines 2’-deoxyctydine and thymidine is somewhat slower than for the purines with apparent \(k_q = 5 \times 10^9 \text{ M}^{-1} \text{ s}^{-1}\). These results are not unexpected since, according to the Rehm–Weller equation,\textsuperscript{37} the excited singlet state of TOTA\textsuperscript{+} is a sufficiently strong oxidant to initiate an electron-transfer reaction at the diffusion-limited rate to form the TOTA radical and a purine radical cation.

The effect of CT-DNA on the fluorescence of TOTA\textsuperscript{+} was examined to assess the effect of binding on the electron-transfer reaction; see Figure 4. As expected, addition of CT-DNA to a TOTA\textsuperscript{+} solution results in efficient fluorescence quenching.


**Figure 2.** UV–Visible absorption spectra for TOTA\textsuperscript{+} in buffer solution (solid line) and in the presence of CT-DNA (dashed line).

**Figure 3.** Induced circular dichroic (CD) spectrum of TOTA\textsuperscript{+} (30 \(\mu\text{M}\) in 10 mM sodium phosphate, pH 7) and sonicated CT-DNA (40 \(\mu\text{M}\), base pairs) recorded in a 10 cm path length cylindrical quartz cell.
are in equilibrium with the same concentration of TOTA. Since all of the DNA samples were determined by extraction into a detergent solution and measurement of its optical absorbance. The apparent affinity of TOTA indicates that dynamic and static quenching processes are result in the rapid injection of a radical cation into the duplex strand.

Interestingly, TOTA is sequence-selective with a preference for C-G base pairs. From the results shown in Figure 5, it is clear that TOTA shows a strong inclination for binding to duplex DNA, and is sequence-selective with a preference for C-G base pairs. Interestingly, TOTA also shows strong binding to triplex DNA. The apparent affinity of TOTA for left-handed Z-DNA increases the length of DNA, resulting in an increased viscosity, which suggests that it binds to DNA by intercalation. However, to RNA. These findings are consistent with the binding of TOTA to duplex DNA by intercalation.

The data of Figure 5 can be used to quantitatively evaluate binding. First, the apparent binding constant may be calculated from the equation \( K_{app} = C_t/C_b (C_T - C_b) \), where \( C_b \) is the amount of TOTA bound plotted in Figure 5, \( C_t \) is the free TOTA concentration (fixed at 1 \( \mu \)M), and \( C_T \) is the total nucleic acid concentration (fixed at 75 \( \mu \)M). For binding to CT-DNA, the data shown in Figure 5 yield \( K_{app} = 4.0 \times 10^4 \) M\(^{-1}\), a value in excellent agreement with the binding constant determined by spectrophotometric titrations as described above. Second, the data for Clostridium perfringens and M. lysodeikticus DNA samples may be used to calculate the ratio \( \alpha = C_b(M. \text{lysodeikticus})/C_b(C. \text{perfringens}) = 3.8 \). The magnitude of \( \alpha \) provides an indication of the nature of the preferred DNA binding site. The value 3.8 corresponds most closely to the preferred binding site composed of adjacent GC base pairs. Photocleavage data, presented later, is fully consistent with such a preferred binding site.

Solution viscosity measurements provide a means to distinguish intercalation from groove binding. In fact, hydrodynamic determination of length changes is among the most stringent tests of the binding mode of ligands to DNA. If a compound binds in a DNA groove, without intercalating, only modest changes in viscosity are generally observed, since this has little effect on the effective length of the polymer. Intercalation, on the other hand, proceeds by unwinding the double helix to accommodate a compound that becomes inserted and stacked between the base pairs. This process results in an effective increase in the DNA contour length. Figure 6 shows the results obtained from viscosity measurements of sonicated CT-DNA (on average 800 base pairs long, \( 2.4 \times 10^{-4} \) M in base pairs) with increasing amounts of bound TOTA\(^+\). Evidently, TOTA\(^+\) increases the length of DNA, resulting in an increased viscosity, which suggests that it binds to DNA by intercalation. However,

the apparent length increase of 1.3 Å per intercalated TOTA\textsuperscript{+} is somewhat less than is observed for other intercalating agents.\textsuperscript{44–46} The length of the DNA samples used in these experiments is longer than the DNA persistence length (≈150 bp), which suggests that the samples will have some flexibility. The lower than expected value for the lengthening could result from changes in the flexibility superimposed on the lengthening.

(2) Cleavage of DNA Photoinduced by TOTA\textsuperscript{+}. The one-electron oxidation of DNA results in the formation of a base radical cation that migrates through the DNA duplex by a hopping mechanism and causes the selective reaction at the 5'-G of GG steps. Reaction at guanine is typically revealed as strand cleavage by treatment of the irradiated sample with hot piperidine. Strand cleavage is detected by gel electrophoresis and autoradiography.\textsuperscript{47} We examined the light-induced reaction of TOTA\textsuperscript{+} bound to the duplex oligonucleotide d(5'-AAAGTGTAACGCG-3'), which was radiolabeled at the A of its 5'-terminus with \textsuperscript{32}P.

Association of TOTA\textsuperscript{+} with the duplex oligonucleotide was confirmed by optical spectroscopy and melting temperature (T\textsubscript{m}) experiments. The T\textsubscript{m} of a 5 μM solution of the duplex oligomer in 10 mM phosphate buffer at pH = 7.0 increases from 52.8 to 57.3 °C when 5 μM TOTA\textsuperscript{+} is added.

Solutions containing the radiolabeled duplex oligomer (5 μM) and bound TOTA\textsuperscript{+} (5 μM; the concentration of free TOTA\textsuperscript{+} in solution is calculated to be ca. 30 nM) were irradiated at 350 nm for 2 h in a Rayonet photoreactor. After irradiation, the DNA was precipitated, treated with hot piperidine, analyzed by denaturing gel electrophoresis, and visualized by autoradiography. Strand cleavage is observed selectively at the 5'-G of GG step of the oligomer; see Figure 7. No cleavage is observed for unirradiated samples. These results indicate that, when bound to DNA, the excited state of TOTA\textsuperscript{+} oxidizes the DNA. This reaction is relatively inefficient, as indicated by the long irradiation time required, which is probably a consequence of its origination from a singlet excited state of TOTA\textsuperscript{+} that generates a singlet radical ion pair that is prone to rapid annihilation by back electron transfer.\textsuperscript{48}

(II) Three-Dimensional Structure of a [d(CGATCG)]\textsubscript{2} TOTA\textsuperscript{2+} Complex. The X-ray data reveal a complex composed of the hexameric duplex plus two TOTA\textsuperscript{2+} molecules intercalated at each terminal base pair step. The DNA residues are numbered 5'-C(1)G(2)A(3)T(4)C(5)G(6)-3' and 5'-C(7)G(8)A(9)T(10)C(11)G(12)-3'. TOTA\textsuperscript{2+} is intercalated between base pairs C(1)-G(12) and G(2)-C(11), and TOTA\textsuperscript{2+}(II) is intercalated between base pairs C(5)-G(8) and G(6)-C(7). Each crystallographic asymmetric unit contains two DNA strands and two TOTA\textsuperscript{2+} molecules, and each [d(CGATCG)]\textsubscript{2} TOTA\textsuperscript{2+} complex contains a noncrystallographic pseudo-2-fold axis between base pairs A(3)-T(10) and T(4)-A(9).

(I) DNA–TOTA\textsuperscript{2+} Interactions. The aromatic faces of the TOTA\textsuperscript{2+} ring systems are engaged in stacking interactions with the adjacent cytosines and guanines. TOTA\textsuperscript{2+} (I) is more centrally disposed relative to the major and minor grooves and seems to engage in tighter interactions with the DNA than TOTA\textsuperscript{2+} (II); see Figure 8. TOTA\textsuperscript{2+} (I) engages in a greater number of contacts (a contact is defined as an interatomic distance < 3.4 Å) with the DNA than TOTA\textsuperscript{2+} (II). TOTA\textsuperscript{2+} (II) is partially extruded out into the major groove. TOTA\textsuperscript{2+} (I) is in contact with nine, and TOTA\textsuperscript{2+} (II) is in contact with seven, atoms of the flanking base pairs. These stacking contacts are significantly longer than some of the TOTA\textsuperscript{2+}–backbone contacts (see below). The formal

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The minor-groove width of the \([d(CGATCG)]_2\):TOTA\(^+\) complex in the central region of the complex is approximately 9 Å, considerably wider than the 7.5 Å width of the \([d(CGATCG)]_2\):daunomycin complex (NDB ID DDF018). The \([d(CGATCG)]_2\):TOTA\(^+\)\(^2\) complex is also bent to a significant extent as compared with DNA/intercalator complexes of similar sequence. When compared with the relatively straight \([d(CGATCG)]_2\):D232 complex, the curvature is quite noticeable. The overall bend angle of the TOTA\(^+\) complex is approximately 18°; the complex is bent toward the major groove.

(3) Water/Cation Region. The solvent region surrounding the DNA:TOTA\(^+\) complex is not particularly well characterized in this structure because the X-ray data were collected at room temperature. However, the electron density maps clearly reveal a magnesium ion with six coordinating water molecules within the major groove of the DNA:TOTA\(^+\) complex. The first-shell water molecules interact with the N7 and O6 positions of G(2) [N7−W2, 2.7 Å; O6−W5, 2.7 Å] and the O4 of T(4) (O4−W3, 2.8 Å). The water molecules of the Mg(H\(_2\)O)\(_5\)\(^+\) are involved in lattice interactions with phosphate groups of a symmetrically related complex. The Mg(H\(_2\)O)\(_5\)\(^+\) is not within contact distance of either TOTA\(^+\) molecule. Seventeen additional water molecules were identified in the electron density map. None of these water molecules is within contact distance with any of the oxygen atoms of TOTA\(^+\). The minor groove of \([d(CGATCG)]_2\):TOTA\(^+\) is devoid of water molecules. Entry of water molecules into the minor groove is precluded by symmetry-related DNA:TOTA\(^+\) complexes. The packing arrangement is similar to that in the \([d(CGATCG)]_2\):D232 complex. A 6-fold screw axis generates duplexes that repeat in the direction parallel to the long axis of each unit cell. Each duplex is tilted such that there is no continuous helix.

(4) Features of the Structure of the \([d(CGATCG)]_2\):TOTA\(^+\)\(^2\) Complex. The essential DNA deformation required for intercalation involves translation of the base pairs in opposing directions along the helical axis, to form an intercalative cavity.\(^{31}\) This base-pair translation disrupts base–base stacking interactions, which are offset at least in part by base–intercalator interactions. The base-pair translations of the TOTA\(^+\) complex are greater than those of other intercalators, and the complex appears to be extended at both intercalation steps. The unusual extension is related to unfavorable electronic or electrostatic interactions between TOTA\(^+\) and DNA; neither of the formal carbocations of the two TOTA\(^+\) molecules is in contact with DNA atoms (<3.6 Å).

DNA deformation can extend along the helical axis into sequences flanking the intercalation site. In the TOTA\(^+\) complex, an unusual deformation is observed at the base pair steps adjacent to the intercalation steps [i.e., at the G(2)−C(11)/A93−T(10) step and at the T(4)−A(9)/C(5)/G(8) steps] inordinately large. The average for the comparison set is 3.18 Å, with the most extreme value of 3.49 Å. The corresponding values in the TOTA\(^+\) complex are 3.60 for TOTA\(^+\)\(^2\) and 3.68 Å for TOTA\(^+\). The unusual helical rise value combine to give the total length of the TOTA\(^+\) complex of 24.7 Å, which is significantly greater than the average of the comparison set (23.6 Å) and is greater than the limiting value (23.7 Å).

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the major groove of the DNA·TOTA\(^+\) complex. The first-shell water molecules interact with the N7 and O6 positions of G(2) and the O4 of T(4). As proposed previously, divalent cations, usually hydrated, and monovalent cations, usually dehydrated, localize adjacent to B-DNA by favorable interaction with the O6 and N7 groups of guanine within the major groove and with the O2 groups of thymine and the N3 groups adenine within the minor groove.\(^{54–57}\) The magnesium ion here follows that pattern of interaction and appears to facilitate the observed helical bend toward the major groove. Localized cations, within G-tract major grooves and A-tract minor grooves, are proposed to interact through electrostatic interactions with DNA phosphate groups and facilitate DNA bending and other deformations (electrostatic collapse). However, as with other X-ray derived structures, the electrostatic interactions within the present complex cannot be fully characterized by the final refined model. With a total anionic charge of 10 and a total cationic charge of 4, the structure here does not achieve charge neutrality. Additional cations, either ordered or disordered, are contained within the solvent region. In some cases lattice effects are known to cause DNA bending in crystals.\(^{58}\) A comparison of the [d(CGATCG)]\(_2\)·TOTA\(^+\) complex with a [d(CGTACG)]\(_2\)·D232 complex suggests that bending of [d(CGATCG)]\(_2\)·TOTA\(^+\) does not originate from lattice interactions. These two complexes have nearly equivalent packing arrangements, including the locations of symmetry-related nucleotides across the minor groove, which appear to widen the minor groove. However, the [d(CGTACG)]\(_2\)·D232 complex is linear.

In summary, these findings show that TOTA\(^+\)(I) and TOTA\(^+\)(II) form crystallographically distinct intercalative complexes that have differing orientations and interactions. TOTA\(^+\)(I) is centrally disposed relative to the major and minor grooves, engages in tighter interactions with the DNA, and is more firmly fixed in position than TOTA\(^+\)(II), which is partially extruded out into the major groove. The shortest DNA·TOTA\(^+\) contacts are not with the flanking bases but with backbone atoms. Electron-deficient and relatively acidic CH groups of TOTA\(^+\)(I) form hydrogen bonds with O4\(^-\) oxygen atoms embedded in the walls of the intercalation cavities. These intracavity C–H\(\cdots\)O hydrogen bonds appear to confer stability and orientation specificity to TOTA\(^+\). TOTA\(^+\)(I) forms two short C–H\(\cdots\)O hydrogen bonds, while TOTA\(^+\)(II) forms one long C–H\(\cdots\)O hydrogen bond. Intracavity hydrogen bonding has not, to our knowledge, been observed previously and may represent a hydrogen bond. Intracavity hydrogen bonding has not, to our knowledge, been observed previously and may represent a hydrogen bond. Intracavity hydrogen bonding has not, to our knowledge, been observed previously and may represent a hydrogen bond. Intracavity hydrogen bonding has not, to our knowledge, been observed previously and may represent a hydrogen bond. Intracavity hydrogen bonding has not, to our knowledge, been observed previously and may represent a hydrogen bond.
of the DNA and the intercalators) was optimized by an energy-gradient minimization method. Subsequently, a histogram of the molecular hydration energies was constructed, and it was employed (together with visual inspection and analysis of the connectivity and coordination of the hydrogen-bond network in the solvation shell) in deciding which water molecules may be discarded, while the more tightly bound ones are maintained. Successive iterations of the above procedure (with each cycle resulting in a smaller number of water molecules retained in the “intimate” hydration environment), culminated in a “minimal hydration shell” (consisting of 120 water molecules) that solvates the intercalated system (including the sodium counterions). The optimized configuration of the intercalated system (with the non-hydrogen atoms of the DNA and the TOTA + intercalators taken from the crystallographical data) is shown in Figure 9, where the counterions and the hydration shell have been excluded for clarity.

To enhance the feasibility of electronic structure calculations on the intercalated system (see below), we have taken advantage of the self-complementary character of the DNA oligomer used in this study. Accordingly, the hydrated six-base-pair intercalated system described above was split into two halves, with each half containing three base pairs, an intercalated TOTA + (I or II), three sodium counterions, and a hydration shell consisting of the 64 water molecules closest to the selected half. The two “half-configurations” were then further annealed (through a cycle of gradual heating and subsequent gradual cooling), resulting in optimal configurations suitable for the quantum calculations described below; atomic configurations of the two hydrated halves are displayed in Figure 13.

(2) Quantum Mechanical Calculations. Figure 10 shows the electronic energy level schemes (only eigenvalues with energies higher than −12 eV are included), calculated for the following: (i) the “stretched” three-base-pair d(5′-C/GA-3′)[4Na] duplex segment of the intercalated hexameric duplex but without the inserted TOTA + (the symbol // represents the expanded interbase distance at the intercalation site, along the axis of the double helix), in which the positions of the DNA atoms were taken from the crystallographic data determined for the intercalated system (except for the hydrogen atoms); (ii) the isolated TOTA +, with crystallographically determined heavy atom positions and with the hydration shell; (iii) the unhydrated TOTA +-intercalated neutral three-base-pair duplex d(5′-C/TOTA +/GA-3′)[3Na], with the DNA and TOTA + coordinates crystallographically determined for the half of the intercalated six-base-pair DNA oligomer that contains TOTA + (I); (iv) same as in column iii but with a hydration shell; and (v) same as in column iv but for the half of the six-base-pair oligomer containing TOTA + (II). Each of the added Na atoms “donates” an electron to the DNA (with the negative charges localizing mostly on the phosphate groups) and thus become the sodium counterions.

Inspection of the eigenvalue schemes of the unhydrated “stretched” three-base-pair DNA duplex segment, d(5′-C/GA-3′)[4Na] [corresponding to the TOTA + (I) site; see Figure 9] and of the isolated TOTA + (columns i and ii) shows that both the highest occupied molecular orbital (HOMO) T and the lowest unoccupied one (LUMO) T of TOTA + lie well below the (HOMO) D of the “stretched” DNA segment, where subscripts T and D indicate the orbital’s predominant association with the TOTA + or the DNA, respectively, where such a designation is meaningful. In the isolated TOTA +, the calculated energy gap between the (HOMO) T and (LUMO) T levels is $\Delta E = 2.08$ eV (to facilitate discussion, we interchangeably refer to $\Delta E$ of the isolated TOTA + also as $\Delta E_T$; see Figure 10). This suggests that charge transfer to the intercalated TOTA + may occur. In this context we note that the energy levels of the unhydrated native three-base-pair DNA segment (not shown) are higher than those of the corresponding unhydrated “stretched” DNA segment (see column i in Figure 10), reflecting the relative instability of the former that is associated with the lack of hydration.

Comparison of the energy level schemes shown in Figure 10 for the unhydrated (iii) and hydrated (iv) intercalated complexes corresponding to the TOTA + (I) site illustrates the marked effect of hydration on the electronic structure of the intercalated DNA system. For the unhydrated complex (iii), the (LUMO) T is almost degenerate with the (HOMO) D (a separation of $\Delta E = 0.06$ eV), with a gap $\Delta = 1.19$ eV between the (LUMO) T and the next higher energy unoccupied molecular orbital, which is associated primarily with the DNA. However, as a result of hydration, the entire electronic energy level spectrum of the

![Figure 10](image_url). Calculated electronic energy level schemes for various components of the intercalation complex. Only eigenvalues with energies higher than −12 eV are displayed. Long horizontal lines correspond to occupied levels, and short lines denote unoccupied ones. The schemes correspond to the following systems (from left to right): (i) “stretched” 3-bp duplex segment d(5′-C/GA-3′)[4Na], where // denotes the empty site; (ii) isolated TOTA +, with the electronic structure calculated for the crystallographically determined heavy atom positions and with the coordinates of the hydrogens determined through optimization employing the Amber 6 force field supplemented by partial charges that were determined as described in the Supporting Information. $\Delta E_T$ denotes the energy gap between the (HOMO) T and (LUMO) T levels that are colored red and blue, respectively. The unoccupied level above the (LUMO) T, i.e., (LUMO) D + 1, is colored green. The level schemes in columns iii and iv correspond to the 3-bp segment shown in column i but with TOTA + inserted at the CG step. The level scheme in column iii was calculated for the unhydrated system (but including the counterions), and the scheme in column iv was calculated for the hydrated system, illustrating the stabilizing effect due to hydration. The colored levels denote the positions of the corresponding eigenvalues of the isolated TOTA + (as described for column ii) but after intercalation. In column iv, $\Delta E_T$ denotes the energy gap between the HOMO and LUMO levels of the complex, and $\Delta E_T$ corresponds to the energy separation between the LUMO and (LUMO) T + 1 levels of the complex. (v) Energy level scheme for the hydrated intercalated complex, calculated for the 3-bp DNA segment associated with TOTA + (II). $\Delta E_T$ denotes the energy separation between the level corresponding to the HOMO of the isolated TOTA + (colored red) and the LUMO of the intercalation complex (colored blue).
complex shifts to lower energies, and the separation between the (HOMO)D and (LUMO)T levels increases significantly (the 3′G and the LUMO level is fully localized on the inserted TOTA\(^+\). Note the correspondence between the (shifted) LDOS of the isolated TOTA\(^+\) and that corresponding to the intercalated TOTA\(^+\)).

Further insights into the nature of the intercalation complex are obtained from analysis of the weighted local density of states (LDOS), where for a given energy eigenvalue the (weighted) probability distribution of the corresponding (normalized) wave function (molecular orbital) in different (spatial) regions of the intercalated system is determined through integration of the square of the wave function in the various regions.\(^{59}\) Plots of the weighted LDOS calculated for the TOTA\(^+\) and TOTA\(^+\)-II segments are shown in Figures 11 and 12, respectively. In these figures, the electronic energy spectrum of the isolated TOTA\(^+\) is uniformly shifted to higher energy so that the (LUMO)\(_T\) of the isolated TOTA\(^+\) will align with the (LUMO)\(_T\) of the intercalated TOTA\(^+\) (shifts of 3.23 and 3.16 eV were used in Figures 11 and 12, respectively). Interestingly, we find that for both intercalation sites the (LUMO)\(_T\) level of the complex is fully localized (that is, we find an LDOS with a weight of unity) on the inserted TOTA\(^+\). Moreover, the (LUMO)\(_T\) calculated for the TOTA\(^+\) and TOTA\(^+\)-II segments are essentially identical to the (LUMO)\(_T\) of the isolated TOTA\(^+\) (see portraits of the orbitals in the top row of Figure 13). These observations, coupled with the finding that the integrated charge in the intercalation region is close to +1, show that intercalation of TOTA\(^+\) into the DNA is not accompanied by transfer of charge from the DNA.

Appendant also from Figures 11 and 12 is that for both intercalation sites the (HOMO)\(_D\) is localized almost entirely on the 3′-G (see also orbital plots in the second row of Figure 13). The (HOMO)\(_T\) and (HOMO 1)\(_T\) states of the isolated TOTA\(^+\) are quasidegenerate (see the red-colored eigenvalues in the column marked isolated TOTA\(^+\) in Figure 10, the top panel of Figure 11, and the corresponding orbital plots in Figure 13). Analysis of the orbital overlaps between these states and various states near the top of the energy spectrum of the intercalation complex allowed us to identify the eigenfunctions of the complex with maximal overlap to the (HOMO)\(_T\) and (HOMO 1)\(_T\) states of the isolated TOTA\(^+\). The states of the complex with maximal overlap are colored red in columns iv and v of

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**Figure 11.** Local densities of states (LDOS) calculated for the isolated TOTA\(^+\) (top panel) and for the hydrated TOTA\(^+\)-II 3-bp DNA segment. The LDOS of the isolated TOTA\(^+\) has been rigidly shifted by 3.23 eV. In each panel we depict the weighted local DOS in the denoted region (see lettered designation in each panel). The height of a vertical bar denotes the weight of the LDOS calculated for the TOTA\(^+\) and that corresponding to the intercalated TOTA\(^+\). Note the correspondence between the (shifted) LDOS of the isolated TOTA\(^+\) and that corresponding to the intercalated TOTA\(^+\). (See portraits of the orbitals in the top row of Figure 13.)

**Figure 12.** Local densities of states, as in Figure 11, but for the hydrated TOTA\(^+\)-II 3-bp duplex DNA segment.

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\(^{59}\) A region (for the purpose of calculating the local density-of-states as well as other local properties) is defined as the union of spheres centered on the atoms that belong to that region, for example, the atoms of a nuclear base, of the sugar-phosphate backbone, of or the TOTA\(^+\) molecule. The radii of the spheres are taken as 2.1 Å for C, N, O, P, and Na and 0.53 Å for a hydrogen atom. Note that with this definition there is some overlap of regions; in particular, the base regions overlap with the backbone region and to a lesser extent with the water region. There is little overlap between the regions of the different bases (the base-to-base distance along the axis of the helix is 3.4 Å) or between a base region and the TOTA\(^+\) region. With this definition of regions we find that the HOMO and LUMO wave functions (see Figures 3 and 4) are found to have a weight in excess of 0.97 in the 3′G and TOTA\(^+\) regions, respectively, with less than 0.02 in any other region.

Figure 13. Orbital portraits for the isolated TOTA\(^+\) (middle column), for the hydrated TOTA\(^+\)(I) 3-bp duplex DNA segments (left column), and for the TOTA\(^+\)(II) segment (right column). Orbital isosurfaces of different signs are distinguished by colors (blue and green), with the contours between the two colors describing the nodal surfaces. The orbital isosurfaces are superimposed on the atomic structures of the corresponding species. The identities of orbitals displayed in the figure are denoted next to each orbital portrait, along with the energy eigenvalue.

Figure 10, and the corresponding orbital portraits are displayed in the third and fourth rows of Figure 13. Interestingly the \((\text{HOMO})_T - (\text{LUMO})_T\) energy gap in the isolated TOTA\(^+\) (\(\Delta_T = 2.08\) eV) is close to the energy separation between the LUMO of the intercalation complex and states of the complex that have maximal overlap with the \((\text{HOMO})_T\) and \((\text{HOMO} - 1)_T\) of the isolated TOTA\(^+\); these energy separations (denoted by \(\Delta_T'\); see Figure 10) are \(\Delta_T' = 2.15\) and 1.95 eV for the TOTA\(^+\)(I) (column iv) and TOTA\(^+\)(II) (column v) intercalation sites, respectively.

These observations pertaining to “spectral rigidity” of the intercalator, together with the aforementioned lack of appreciable electron transfer to TOTA\(^+\), lead us to conclude that electrostatic forces and induced polarization of the electronic distribution underlie the intercalation and binding process of TOTA\(^+\) into DNA. This is further supported by the polarization charge plots shown in Figure 14, where the polarization charge is calculated as the charge difference: 
\[
\delta \rho = \rho\{\text{d}(5'-C//\text{GA}-3')\} - \rho\{\text{d}(5'-C//\text{GA}-3')\} + \rho\{\text{Na}\} - \rho\{\text{d}(5'-C//\text{GA}-3')\} + \rho\{\text{Na}\}
\]
where the species in the first term is neutral, that in the second term is singly negatively charged, and the third is TOTA\(^+\). The upper panel shows \(\delta \rho\) for the entire intercalated three-base-pair DNA segment [TOTA\(^+\)(I)], the plot shown in the middle panel excludes the intercalator region [5], and in the plot shown in the bottom of the panel the charge differences in the regions of the intercalator and the hydrating shell are excluded. On the left side of the middle panel, we display the charge difference in the TOTA\(^+\) region extracted from the complex, viewed from above the plane of the molecule. In these charge-difference portraits, magenta regions correspond to those where excess electronic charge is induced by the intercalation process, and yellow regions correspond to those where depletion of the electronic charge is found. The induced polarization charge on the inner surface of the “intercalation cavity” (defined by the 5′-3′G and G-C base pairs surrounding the inserted TOTA\(^+\) from above and from below, respectively) is \(-0.36\)e (obtained by integrating over the magenta region of the cavity). The integrated depletion in the adjacent (yellow) region is 0.25e (the total polarization charge, integrated for the entire system, vanishes).

The interaction of the TOTA\(^+\) with the negatively charged hydrated duplex \(\text{d}(5'C//\text{GA}-3')[\text{Na}]\), as well as with the polarization charge discussed above, accounts for the near rigid shift of the TOTA\(^+\) region in the electronic energy-level spectrum relative to the spectrum of the isolated TOTA\(^+\). The effect of hydration is to lower the occupied levels of the DNA component of the spectrum, which opens the gap \(\Delta_p\) between the \((\text{HOMO})_T\) of the complex (localized on the 3'-G) and the \((\text{LUMO})_T\) of the complex (localized on the intercalated TOTA\(^+\)). The water molecules have little effect on the TOTA\(^+\) part of the spectrum because their affinity toward the TOTA\(^+\) is relatively small compared with their binding affinity to the DNA bases.

Conclusions

The experimental and computational studies reported here present a clear picture of the photochemical, structural, and electronic consequences of the interaction of TOTA\(^+\) with duplex DNA. Intercalation is a complex phenomenon that distorts the DNA structure and causes electronic energy level shifts.

Solution-phase experiments indicate that TOTA\(^+\) binds to DNA with a preference for G-C base pairs and triplex structures. The X-ray crystallographic data show unambiguously that
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The X-ray crystallographic data reveal that intercalation of TOTA+ results in an unusually large extension of the helical rise of the DNA. This seems to be a result of unfavorable electronic interactions between TOTA+ and the bases that form the intercalation site. The orientation of TOTA+ is sensitive to hydrogen-bonding interactions with backbone atoms. In particular, intercavity hydrogen bonds provide stability and orientational specificity.

The electronic structure calculations reveal the importance of backbone, water, and counterions, which shift the energy levels of the DNA bases and the TOTA+ orbitals significantly and differently. Importantly, the calculations show no meaningful charge transfer from DNA to TOTA+ because (LUMO)T of TOTA+ falls in the gap of (HOMO)D and (LUMO)D. Consequently, the important electronic interactions are repulsive between occupied orbitals of TOTA+ and occupied orbitals of the DNA bases. It is clear from the X-ray data and from the electronic calculations that electrostatic interactions control the intercalation of TOTA+ in DNA. TOTA+ strongly polarizes the intercalation cavity, where a sheet of excess electron density is induced that surrounds the TOTA+.

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Supporting Information Available: Two tables listing crystallographic and refinement statistics for d(CGATCG)2·TOTA+2 and base-step distances and overall DNA length of intercalated hexamers, diagram and calculation of partial charges on the TOTA cation, and preparation of hydrated samples for calculation (PDF). This information is available free of charge via the Internet at http://pubs.acs.org.

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